Effects of growth temperature on the photovoltaic properties of RF sputtered undoped NiO thin films

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ARTICLE INFO

Keywords: Magnetron sputtering Material properties NiO thin film Substrate temperature Transparent conducting oxide

ABSTRACT

In this study, nickel oxide (NiO) thin films were deposited on soda lime glass using radio-frequency magnetron sputtering at different growth (substrate) temperatures ranging from room temperature (RT) to 400 °C. The effects of substrate temperature on the structural, morphological, electrical, and optical properties were investigated. The XRD pattern unveiled a dominant peak with (2 0 0) preferential orientations for the film grown at 100 °C. However, for samples grown at high temperatures, a gradual decrease of (2 0 0) peak intensity was observed, which may be the result of the decomposition of NiO as confirmed via EDX. Surface morphology from FESEM revealed that grains were randomly orientated on the surface with maximum grain size of 19.43 nm. Upon increasing the growth temperature, the crystal quality and grain size substantially deteriorated, which is consistent with the XRD results. Scanning probe microscopy (SPM) finds rough surface with the highest surface roughness obtained at RT with a value of 1.232 nm. Electrical resistivity was found to be highly dependent on the growth temperature, the crystal quality and grain size substantially deteriorated, which is consistent with the XRD results. Scanning probe microscopy (SPM) finds rough surface with the highest surface roughness obtained at RT with a value of 1.232 nm. Electrical resistivity was found to be highly dependent on the growth temperature that decreases from 2150 \(\Omega\) cm to 72 \(\Omega\) cm as the substrate temperature increases. For optical properties, the optical bandgap of the NiO films decreases from 3.8 eV to 3.2 eV as a function of substrate temperature as derived from the optical transmittance data. Results show the potential application of the NiO films in photovoltaic devices.

Introduction

Nickel oxide (NiO) is a semi-transparent and naturally p-type semiconductor, which has a wide bandgap ranging from 3.6 eV to 4.0 eV. Recently, NiO has drawn tremendous attention owing to its exceptional optical, electrical, and magnetic properties and good chemical stability. The remarkable inherent properties of NiO has enabled this compound to become a promising material for photovoltaic solar cells [1], electrochromic devices [2], ultraviolet photo-detectors [3], and gas sensors [4]. In solar cells, two main applications are anticipated for NiO: either as a transparent conducting oxide (TCO) [5] or hole transporting layer (HTL) in perovskite solar cells [6,7]. In particular, replacing organic HTL to inorganic is important because the latter has high hole mobility, easily processable properties, low production cost, and highly stable [8]. Deposition temperature is another crucial issue to grow any layer on the top of the perovskite absorber because it easily decomposes when exposed at > 150 °C, although the precise value is disputable [9]. Several physical and chemical methods have been used to fabricate NiO film, including spray pyrolysis [10], radio frequency (RF) magnetron sputtering [11,12], direct-current (DC) magnetron
Experimental procedure

NiO thin films were deposited via RF magnetron sputtering onto well-cleaned soda lime glass (SLG) substrates (30 mm × 30 mm × 1.1 mm) at room temperature. Substrates were cleaned ultrasonically in soap water–deionized water–methanol–isopropanol–deionized water sequence. Thereafter, the substrates were dried in N₂ gas flow and rinsed in acetone. Lastly, substrates were exposed under UV for 30 min to remove organic residues. Stoichiometric NiO sputtering target (diameter: 2 in.; purity 99.99%, custom-made) was used as the source material. The diameter and thickness of the NiO sputtering target were 50 mm and 5 mm, respectively. The distance between target and substrate was maintained at 14 cm. Prior to sputtering, the target was pre-sputtered for 15 min at 50 W power to clean the target surface. The base pressure, working pressure, sputtering power, and deposition time were maintained at 5 × 10⁻⁶ Torr, 28 m Torr, 50 W, and 90 min, respectively. The reactive gas was high-purity Ar (99.9999%) and the flow rate was 4 sccm as controlled by a mass flow controller. The rotation speed of the substrate holder was maintained at 10 rpm during sputtering. Prior to sputtering, the substrates were preheated for 45 min at intended deposition temperature of 100–400 °C.

BRUKER axs-D8 Advance Cu-Kα X-ray diffractometer was used to measure the structural and crystallographic properties of NiO thin films. X-ray diffraction (XRD) patterns were recorded in the diffraction angle (2θ) ranging from 20° to 70° using Cu Kα radiation wavelength, λ = 1.5406 Å. NT–MDT, NTEGRA PRIMA scanning probe microscopy (SPM) in noncontact mode setting was used to investigate the surface topography and roughness of the NiO thin film. Optical transmission and absorbance measurements were performed using Perkin Elmer Lambda 950 UV/Vis/NIR spectrophotometer in the wavelength range of 200–1000 nm. The Hall Effect Measurement system called HMS ECOPIA 3000 (magnetic field 0.57 T, probe current 10 mA) was used to measure the electrical parameters, including carrier concentration, mobility, and resistivity. Lastly, SUPRA 55 VP field emission scanning electron microscope (FESEM) was used to measure the film thickness, grain size, surface morphologies, and cross-sectional view with accelerating voltage of 10–20 kV, which is equipped with energy dispersive x-ray (EDX) detectors.

Results and discussion

Structural analysis of the NiO films

The XRD pattern of the NiO thin films with diffraction angles from 20° to 70° at different substrate temperatures is shown in Fig. 1(a). The diffraction patterns of the RF-sputtered NiO films show that all samples exhibit two diffraction peaks corresponding to NiO (2 0 0) and Ni (2 2 0) orientations at 43.3° and 62.5°, respectively, except the sample sputtered at 400 °C. The XRD pattern from this study is appropriately matched with the standard XRD spectrum documented in JCPDS card No. 00-047-1049, thereby representing that the grown polycrystalline NiO layer crystallizes in a cubic structure. The dominant peak intensity
is seen at substrate temperature 100 °C, thereby indicating strong orientation and good crystallinity. The (2 0 0) intensity of NiO diffraction peak steadily decreases as substrate temperature continues to increase. At 400 °C, the crystal quality of NiO films are entirely degraded and diffraction peaks along (1 1 1) and (1 1 0) planes appeared at 35.89° and 30.92°, respectively. At substrate temperature of over 100 °C, the crystallinity of NiO films is degraded by the high thermal energy available to the adatoms of NiO films. This result may be ascribed to low atomic mobility at low substrate temperature, thereby leading to the formation of a preferred orientation along the (2 0 0) plane. However, when the substrate temperature increased, the NiO atoms gain additional energy and tend to move another preferred orientation along the (1 1 1) and (1 1 0) planes. This experiment was conducted via RF sputtering of the NiO target. Thus, NiO is partially dissociated into metallic Ni and O at high substrate temperature. This phenomenon leads to a reduction of the crystal quality of NiO thin films. These observations are in good agreement with prior studies that reported a decrease in the crystallinity of different thin films owing to their decomposition at high substrate temperatures [11,24]. Several structural parameters, including crystallite size, FWHM, strain, dislocation density, microstrain, and lattice parameter, were calculated from the XRD data. Fig. 1(b) represents the variations of crystallite size and FWHM of the NiO diffraction peaks along the (2 0 0) plane corresponding to different substrate temperatures. The average crystallite size \(d_{hkl}\) was estimated using the Debye–Scherrer’s equation [25]:

\[
d_{hkl} = \frac{K \lambda}{B \cos \theta}
\]

where \(\lambda\) is the X-ray wavelength, \(B\) is the FWHM intensity of the main peak observed at \(2\theta\) in radian, \(\theta\) is the Bragg’s angle of diffraction, and \(K\) is a constant. Fig. 1(b) shows that the crystallite size of NiO decreases with increasing substrate temperature, although FWHM increases. The crystallite size decreased from 19.43 nm to 13.10 nm as the substrate temperature increased from RT to 400 °C. This result is ascribed to the low surface diffusion of the adatoms on the substrate because of the decomposition of the NiO thin film at high substrate temperature [11].

The microstrain (\(\varepsilon\)) and dislocation density (\(\delta\)) of the NiO thin films were calculated using Eqs. (2) and (3), respectively [26,27]:

\[
\varepsilon = \frac{\beta}{(4\tan \theta)}
\]

\[
\delta = \frac{n}{D^2}
\]

where \(D\) is the grain size and \(n\) is a factor that is nearly equal to the unity for minimum dislocation density.

Fig. 2(a) depicts the variations of the calculated dislocation densities and microstrain for the (2 0 0) crystallographic orientation of the NiO films sputtered at different substrate temperatures. The increasing degree of microstrain and dislocation density is found for films sputtered at high substrate temperatures. Grain size enlargement through grain fusion at high substrate temperature reduces the lattice point mismatch of two adjacent crystals. Lattice point mismatch initiates from the misregistry of the lattice in one part of the crystal with respect to another part along the interconnecting grain boundary. Fig. 2(a) shows decreased dislocation density and microstrain at low substrate temperature caused by low lattice point mismatch that eventually becomes favorable for photovoltaic application. The collective impact of the increase in the microstrain and dislocation density can be used to clarify the substantial enhancement in the stacking fault of films with increasing substrate temperature [28].

The variation of the lattice constant and percentages of change in the lattice constant (strain) along the plane (2 0 0) crystallographic orientation as a function of substrate temperature are presented in Fig. 2(b). The lattice constant of the NiO thin films was calculated using Eq. (4) [29]:

\[
\frac{1}{d_{hkl}} = (h^2 + k^2 + l^2)^{1/2}
\]

Fig. 2. (a) Effect of growth temperature on microstrain and dislocation density (b) lattice constant and strain of the NiO thin film.

The dotted line in the graph indicates the equilibrium lattice constant value, which corresponds to 0% strain. The lattice constant and strain of the NiO film increased with substrate temperature. Growing substrate temperature encourages lattice relaxation with film sputtered at 100 °C maintaining equilibrium lattice parameter without strain. Fig. 2(b) reveals that the equilibrium lattice constant of the RF-sputtered NiO at 100 °C is 4.17 Å at zero strain. This result is consistent with the previous experimental finding of Madhavi et al. [17]. The variations in the lattice constant and strain with substrate temperature are due to the tensile stress developed in the films at high substrate temperature [30]. Consequently, the film will crack or shrink because of compressive strain, which could not be observed by the naked eye but is detrimental for photovoltaic application.

Surface morphology of the NiO films

The variation of surface morphology and the cross-sectional view of the NiO thin films at different substrate temperatures were observed via FESEM (see Fig. 3). The FESEM micrograph demonstrates the nanoscale grain of the NiO thin films and the grain size is strongly dependent on the substrate temperature. The FESEM micrograph of NiO deposited at 100 °C possess uniform grain size and high packing density with an average grain of ~16 nm. Grain size and crystal quality rapidly deteriorated with increasing substrate temperature. Interestingly, the grain shapes completely become irregular at 400 °C. This morphological observation is in good agreement with the previously discussed structural analysis of the NiO thin film. From the cross-sectional view, thickness gradually decreases as substrate temperature increases. Thickness is reduced from 100 nm to 60 nm as substrate temperature increased from RT to 100 °C. This outcome may be the result of the dependency
between grain size and film thickness [31]. The effect of Ni/O ratio on the substrate temperature is shown in Fig. 4. The ratio of Ni/O increased with substrate temperature because the NiO atoms decomposed at high substrate temperature, thereby leading to the degradation of crystal quality. This observation reveals a correlation among the structural properties of NiO thin film and is consistent with the findings of Ahmed et al. and Reddy et al. [11,32].

The surface roughness analysis of NiO thin films deposited at

<table>
<thead>
<tr>
<th>Substrate temperature, °C</th>
<th>Carrier concentration, cm$^{-3}$</th>
<th>Carrier mobility, cm$^{2}$/Vs</th>
<th>Resistivity, $\Omega$ cm</th>
</tr>
</thead>
<tbody>
<tr>
<td>RT</td>
<td>$6.08 \times 10^{14}$</td>
<td>6.73</td>
<td>1503</td>
</tr>
<tr>
<td>100</td>
<td>$2.85 \times 10^{14}$</td>
<td>10.20</td>
<td>2150</td>
</tr>
<tr>
<td>200</td>
<td>$4.49 \times 10^{14}$</td>
<td>6.75</td>
<td>2060</td>
</tr>
<tr>
<td>300</td>
<td>$7.32 \times 10^{15}$</td>
<td>2.87</td>
<td>298</td>
</tr>
<tr>
<td>400</td>
<td>$3.87 \times 10^{16}$</td>
<td>2.23</td>
<td>72</td>
</tr>
</tbody>
</table>

Fig. 3. Top and cross-sectional (inset) views of the NiO thin film at different substrate temperatures: (a) RT, (b) 100 °C, (c) 200 °C, (d) 300 °C, and (e) 400 °C.

Fig. 4. Variation in the Ni/O ratio as a function of substrate temperature.

Fig. 5. 3D SPM images of the NiO thin films deposited by RF sputtering at different substrate temperatures: (a) RT, (b) 100 °C, (c) 200 °C, (d) 300 °C, and (e) 400 °C.

Table 1
Hall parameters of the NiO thin films at different substrate temperatures.

between grain size and film thickness [31]. The effect of Ni/O ratio on the substrate temperature is shown in Fig. 4. The ratio of Ni/O increased with substrate temperature because the NiO atoms decomposed at high substrate temperature, thereby leading to the degradation of crystal quality. This observation reveals a correlation among the structural properties of NiO thin film and is consistent with the findings of Ahmed et al. and Reddy et al. [11,32].
different substrates (Si wafer) was carried out through SPM. The 3D SPM images of the NiO thin films are shown in Fig. 5. The average roughness of the substrate is defined as the proportion of the root means square (RMS) value to the average height. The values of the RMS surface roughness at substrate temperatures RT, 100 °C, 200 °C, 300 °C, and 400 °C are observed to be 1.232 nm, 1.046 nm, 0.681 nm, 0.632 nm, and 0.621 nm, respectively. Hence, surface roughness decreased as substrate temperature increased, thereby correlating with the structural analysis via XRD. The crystalline quality of the NiO thin films has been improved at low temperature. This property may help induce additional ordered packing of molecules to eliminate voids and dislocations and can act as charge traps. A similar experimental output is also reported by Fasaki et al. [31].

**Electrical properties of the NiO films**

Table 1 shows the electrical properties of the NiO thin films. The growth of these films at various substrate temperatures unveils a p-type conductivity with carrier concentration from \(2.85 \times 10^{14} \text{cm}^{-3}\) to \(3.87 \times 10^{16} \text{cm}^{-3}\), thereby reflecting the substantial influence of substrate temperature on the electrical properties of NiO. The variation of electrical conductivity of the NiO thin films at different substrate temperatures depends on the microstructural defects of the NiO crystal and surface chemical reaction [23]. The XRD studies indicate that the crystallinity structure decreases as the substrate temperature increases owing to microstructural defect, thereby increasing carrier concentration [33]. The relation among carrier concentration, carrier mobility, and resistivity can be explained by Eqs. (5) and (6):

\[
\sigma = \frac{q\eta\mu_e}{\rho} 
\]

where \(\sigma\), \(\rho\), \(q\), \(\eta\), and \(\mu_e\) are conductivity, resistivity, electrical charge, carrier concentration, and carrier mobility, respectively. From the equations, electrical mobility and carrier concentration are inversely proportional to the resistivity. The carrier mobility of the NiO thin film initially increased with the substrate temperature along with decreasing resistivity. After 100 °C, the carrier mobility gradually decreased with continuous increasing substrate temperature and decreasing resistivity. This outcome could be caused by the increase in the dislocation density as presented in Fig. 2(a), which promotes carrier scattering and subsequently reduces carrier mobility. The NiO thin films showed a high electrical resistivity of 2150 \(\Omega\) cm at a substrate temperature of 100 °C. The resistivity of the NiO thin films decreased to 72 \(\Omega\) cm as substrate temperature increased to 400 °C because of increasing carrier
concentration, which is induced by Ni-rich composition through the partial decomposition of NiO at a high substrate temperature.

**Optical characteristics of the NiO films**

Substrate temperature strongly influences the optical transmittance of samples [24]. The optical transmittance of NiO was measured in the wavelength range of 200–1000 nm, as shown in Fig. 6(a). In the visible range, the transmittance of the NiO thin films decreases moderately as substrate temperature increases. The transmittance of these films is approximately 83.69% but drastically decreases thereafter to 60.56% as the substrate temperature increases from RT to 400 °C. These findings may be attributed to the increase in Ni content in NiO thin films, which coincidences with compositional analysis of this study. Pronounced light scattering is expected because of the large amount of grain boundaries and point defects that reflect the incident light [34].

The optical bandgap (E_g) of the NiO thin films at different substrate temperatures was calculated from the transmission spectra using the following equation [35]:

\[ \alpha h\nu = A(h\nu - E_g) \]

where \( \alpha \) is the absorption coefficient, \( h \) is Planck’s constant, \( \nu \) is the incident photon frequency, \( A \) is the constant, and \( E_g \) is the optical bandgap. Fig. 6(b) shows the calculated \( E_g \) at different substrate temperatures and shows the values ranging from 3.80 eV to 3.20 eV for changing the substrate temperature from RT to 400 °C, which is fitting in the range of standard reported value. \( E_g \) narrows as substrate temperature increases as presented in Fig. 7(a). The change in \( E_g \) is due to the change in the stoichiometry and crystallinity of the NiO thin films [36]. Moreover, \( E_g \) may change because of the shift in absorption edge and changing carrier concentration, which can be explained by the Burstein–Moss effect [32,37]. The Burstein–Moss effect indicates that the Fermi level moves into the conduction band of a degenerated semiconductor that leads to energy band expansion [38]. Thus, the shrinkage effect is leading over the Burstein–Moss effect because \( E_g \) is reduced with the increasing substrate temperature.

The Urbach tail is the exponential region in the absorption spectrum, which is a function of the absorption coefficient and photon energy. It reveals the shifts between the tails of the conduction and valence bands. The Urbach tail appears in the poor-crystalline, disordered, and amorphous materials owing to the localized states of these materials, which are extending or narrowing the bandgap. Rahal et al. also reported that Urbach energy is closely related to the disorder in the materials, which are extending or narrowing the bandgap. Rahal et al. also reported that Urbach energy is closely related to the disorder in the film network [39]. In the low photon energy range, the Urbach empirical rule can be explained by the following equation [40]:

\[ \alpha = \alpha_o \exp(h\nu/E_u) \]

where \( \alpha \) is the absorption coefficient, \( h\nu \) is the photon energy, \( \alpha_o \) is a constant, and \( E_u \) represents the Urbach energy, which is dependent upon the temperature. The Urbach energy of the NiO thin films have been estimated from the slopes of \( \ln(\alpha) \) versus photon energy (\( h\nu \)) plots (see Fig. 7(b)).

The inverted slope of \( \ln(\alpha) \) versus photon energy (\( h\nu \)) plots determines the values of the Urbach energy (see Table 2).

The Urbach energy of the NiO films increases from 177 meV to 438 meV as substrate temperature increases from RT to 400 °C. The highest Urbach energy is found at 400 °C, thereby indicating that the crystallinity becomes inferior and the degree of defect increases with temperature [41]. These defects and disorders may lead to the formation of delocalized state near the band level and the enhancement of the Urbach energy value. This finding offers appreciated information for the research community that the RF-sputtered NiO thin films deposited at low substrate temperature are beneficial for photovoltaic applications. The structural, optical, and electrical properties of the NiO thin films have been compared with previously reported findings (see Table 3).

Hence, the core scientific findings from this experimental study can be deduced as follows. The loss of O from the NiO thin film during high-temperature sputtering leads to Ni-rich composition, which ultimately reduces optical transmission (as well as the optical band gap) and electrical resistivity. Therefore, the careful selection of growth temperature is imperative to tailor the material properties of the NiO thin film according to their intended applications (e.g., TCO, HTL, buffer layer).

**Conclusions**

NiO thin films were deposited via RF magnetron sputtering at different substrate temperatures from RT to 400 °C on the top of a soda lime glass. Leading peak intensity was obtained at 100 °C and the crystal quality substantially deteriorated as the increase of substrate temperature over 100 °C. Surface roughness, film thickness, and crystallite size decreased as the substrate temperature increased. Microstrain and dislocation density were enhanced by the substrate temperature probably due to the released stress in the NiO thin film. The atomic percentage of Ni increased at high substrate temperature for a possible decomposition of O. The carrier concentration and carrier mobility of the NiO thin film increased, whereas resistivity decreased as the substrate temperature increased. The lowest resistivity obtained was 72.0 Ω cm at 400 °C. The optical properties revealed that the optical bandgap shrinks corresponding to the increase of Urbach energy at high substrate temperature. The results prove that the properties of the NiO thin film deposited at 100 °C are suitable for photovoltaic application, particularly as HTM of the perovskite solar cell.

**Acknowledgement**

This study was financially supported by the Fundamental Research Grant Scheme (Code: FRGS/1/2017/TK07/UKM/02/9) of the Ministry of Higher Education, Malaysia. The authors would also like to extend their appreciation to the Deanship of Scientific Research at King Saud University for funding this study through Research Group No. RGP-1438-025.

**Appendix A. Supplementary material**

Supplementary data to this article can be found online at https://doi.org/10.1016/j.rinp.2019.102360.

**References**


